

# The Korteweg-de Vries equation and theta functions

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We consider the Korteweg-de Vries (KdV) partial differential equation

$$4u_t - 6uu_x - u_{xxx} = 0 \tag{1}$$

The function  $u = u(x, t)$  models the height of one-dimensional waves in shallow water.

## 1 Translation waves solutions and plane cubics

We look for translation waves solutions of the form

$$u(x, t) = 2 \cdot v(x + a \cdot t)$$

for a constant  $a$  and a function  $v = v(z)$  that we take to be as nice as possible, meaning analytic. Imposing that this is a solution to the KdV we get

$$8a \cdot v' - 24vv' - 2v''' = 0$$

This is a differential equation for the function  $v$ . Integrating this we get

$$8av - 12v^2 - 2v'' + b = 0$$

where  $b$  is an integration constant. We can multiply this equation by  $v'$  to obtain

$$8avv' - 12v^2v' - 2v''v' + bv' = 0$$

and we can integrate this again to obtain

$$4av^2 - 4v^3 - (v')^2 + bv + c = 0$$

where  $c$  is again an integration constant. In conclusion, we proved the following

**Lemma 1.1.** *Let  $v$  be an analytic function which satisfies a differential equation of the form*

$$(v')^2 = -4 \cdot v^3 + 4a \cdot v^2 + b \cdot v + c$$

*for certain constants  $a, b, c$ . Then the function  $u(x, t) = 2v(x + at)$  is a solution to the KdV equation.*

In geometric terms, this leads naturally to the plane cubic curve

$$E = \{(x, y) \in \mathbb{C}^2 \mid y^2 = -4 \cdot x^3 + 4a \cdot x^2 + b \cdot x + c\}$$

and we are looking for an analytic function  $v$  such that the map

$$(v, v'): \mathbb{C} \longrightarrow E, \quad z \mapsto (v(z), v'(z))$$

ends up into  $E$ . A plane cubic curve in the form  $\{y^2 = f(x)\}$ , where  $f(x)$  is a polynomial of degree three is said to be in Weierstraß form, and Weierstraß actually found a function that satisfies the differential equation of Lemma 1.1. This function is known as the Weierstraß  $\wp$ -function. However, here we will consider it through the point of view of the theta function.

## 2 Complex tori and theta functions

Let  $\mathcal{H} = \{\tau \in \mathbb{C} \mid \text{Im } \tau > 0\}$  be the complex upper-half plane, consisting of complex numbers with positive imaginary part.

**Definition 2.1** (Theta function). The theta function is

$$\theta: \mathbb{C} \times \mathcal{H} \rightarrow \mathbb{C}, \quad \theta(z, \tau) = \sum_{n \in \mathbb{Z}} \exp(\pi i n^2 \tau + 2\pi i n z)$$

**Remark 2.2.** Since we are imposing that  $\text{Im } \tau > 0$ , the series in the definition converges uniformly on compact subsets, so that the theta function is a well-defined holomorphic function.

The most important property of the theta function is its quasiperiodicity:

**Proposition 2.3** (Quasiperiodicity). *Let  $m, n \in \mathbb{Z}$ . Then*

$$\theta(z + m + \tau n, \tau) = \exp(-\pi i n^2 \tau - 2\pi i n z) \cdot \theta(z, \tau)$$

*Proof.* We can compute

$$\begin{aligned} \theta(z + m + \tau n, \tau) &= \sum_{h \in \mathbb{Z}} \exp(\pi i h^2 \tau + 2\pi i h(z + m + n\tau)) \\ &= \sum_{h \in \mathbb{Z}} \exp(\pi i h^2 \tau + 2\pi i h n \tau + 2\pi i h z) \exp(2\pi i m h) \\ &= \sum_{h \in \mathbb{Z}} \exp(\pi i (h + n)^2 \tau - \pi i n^2 \tau + 2\pi i (n + h)z - 2\pi i n z) \\ &= \exp(-\pi i n^2 \tau - 2\pi i n z) \cdot \sum_{h \in \mathbb{Z}} \exp(\pi i (h + n)^2 \tau + 2\pi i (n + h)z) \\ &= \exp(-\pi i n^2 \tau - 2\pi i n z) \cdot \theta(z, \tau) \end{aligned}$$

□

This quasiperiodicity is interpreted geometrically through the action of the lattice  $\Lambda_\tau = \{m + \tau n \mid m, n \in \mathbb{Z}\}$  on the complex plane  $\mathbb{C}$ . One can see that the quotient  $\mathbb{C}/\Lambda_\tau$  is topologically equivalent to a real torus. Since it has also a complex structure inherited from  $\mathbb{C}$  it is also called a complex torus. Before looking at some consequences of the quasiperiodicity, let's see another useful property of the theta function

**Lemma 2.4** (Parity). *The theta function is even, meaning that*

$$\theta(-z, \tau) = \theta(z, \tau)$$

*Proof.* We can simply compute

$$\theta(-z, \tau) = \sum_{n \in \mathbb{Z}} \exp(\pi i n^2 \tau - 2\pi i n z) = \sum_{n \in \mathbb{Z}} \exp(\pi(-n)^2 \tau + 2\pi i(-n)z) = \theta(z, \tau)$$

□

One immediate consequence of the quasiperiodicity is on the logarithmic derivatives of the theta function

**Lemma 2.5.** *For any  $n, m \in \mathbb{Z}$  it holds that*

$$\begin{aligned} \frac{\partial \log \theta}{\partial z}(z + m + \tau n, \tau) &= -2\pi i n + \frac{\partial \log \theta}{\partial z}(z, \tau), \\ \frac{\partial^k \log \theta}{\partial z^k}(z + m + \tau n, \tau) &= \frac{\partial^k \log \theta}{\partial z^k}(z, \tau) \end{aligned} \quad \text{for all } k \geq 2$$

*Proof.* Just take the logarithm in the quasiperiodicity and then the successive derivatives. □

With this we can identify the zeroes of the theta function

**Proposition 2.6.** *For each fixed  $\tau \in \mathcal{H}$  the zeroes of the corresponding theta function are*

$$\{z \mid \theta(z, \tau) = 0\} = \left\{ \frac{1}{2} + \frac{1}{2}\tau \right\} + \Lambda_\tau$$

*furthermore, these are all simple zeroes.*

*Proof.* The quasiperiodicity shows that all the zeroes are  $\Lambda_\tau$  invariant. We are now going to show that, up to translation by  $\Lambda_\tau$ , the theta function has a unique zero, which is moreover a simple zero. To do so, draw a fundamental parallelogram  $P$  for  $\Lambda_\tau$  whose sides do not contain any zero. Then we want to show that there is a unique zero inside this fundamental parallelogram. The Residue Theorem applied to the logarithmic derivative shows that

$$\#\{\text{zeroes of } \theta \text{ inside } P\} = \frac{1}{2\pi i} \int_{\partial P} \frac{\partial \log \theta}{\partial z} dz$$

however, the first relation of Lemma 2.5 shows that the integral on the right is exactly  $2\pi i$ . Hence,  $\theta$  has only one zero inside  $P$ , counted with multiplicity. To conclude, we need to

prove that  $z_0 = \frac{1}{2} + \frac{1}{2}\tau$  is a zero of theta. We use a trick following Jacobi: consider the function

$$\theta_{11}(z, \tau) = \exp\left(\frac{1}{4}\pi i\tau + \pi i\left(z + \frac{1}{2}\right)\right) \cdot \theta\left(z + \frac{1}{2} + \frac{1}{2}\tau, \tau\right)$$

We claim that  $\theta_{11}(z, \tau)$  is odd, meaning that  $\theta_{11}(-z, \tau) = -\theta_{11}(z, \tau)$ . If this is true, then  $\theta_{11}(0, \tau) = 0$ , meaning that  $\theta\left(\frac{1}{2} + \frac{1}{2}\tau, \tau\right) = 0$ . To prove that  $\theta_{11}$  is odd we compute

$$\begin{aligned} \theta_{11}(-z, \tau) &= \exp\left(\frac{1}{4}\pi i\tau + \pi i\left(-z + \frac{1}{2}\right)\right) \cdot \theta\left(-z + \frac{1}{2} + \frac{1}{2}\tau, \tau\right) \\ [\theta \text{ is even}] &= \exp\left(\frac{1}{4}\pi i\tau + \pi i\left(-z + \frac{1}{2}\right)\right) \cdot \theta\left(z - \frac{1}{2} - \frac{1}{2}\tau, \tau\right) \\ &= \exp\left(\frac{1}{4}\pi i\tau + \pi i\left(-z + \frac{1}{2}\right)\right) \cdot \theta\left(z + \frac{1}{2} + \frac{1}{2}\tau - 1 - \tau, \tau\right) \\ [\text{quasiperiodicity}] &= \exp\left(\frac{1}{4}\pi i\tau + \pi i\left(-z + \frac{1}{2}\right)\right) \cdot \exp\left(-\pi i\tau + 2\pi i\left(z + \frac{1}{2} + \frac{1}{2}\tau\right)\right) \\ &\quad \cdot \theta\left(z + \frac{1}{2} + \frac{1}{2}\tau, \tau\right) \\ &= \exp(\pi i) \cdot \exp\left(\frac{1}{4}\pi i\tau + \pi i\left(z + \frac{1}{2}\right)\right) \cdot \theta\left(z + \frac{1}{2} + \frac{1}{2}\tau, \tau\right) \\ &= -\theta_{11}(z, \tau). \end{aligned}$$

Remark: in the previous reasoning we did not need the factor  $\exp\left(\frac{1}{4}\pi i\tau\right)$ , but it is there for another reason (that we do not care about today) and we kept it for notational consistence.  $\square$

As an immediate consequence of this, we get

**Corollary 2.7.** The meromorphic functions  $\frac{\partial^2 \log \theta}{\partial z^2}, \frac{\partial^3 \theta}{\partial z^3}, \left(\frac{\partial^2 \log \theta}{\partial z^2}\right)^2$  are  $\Lambda_\tau$  periodic, and, up to translations by  $\Lambda_\tau$  they have a unique pole at  $z_0 = \frac{1}{2} + \frac{1}{2}\tau$ . Furthermore, this point is a pole of order 2, 3, 4 respectively.

*Proof.* If  $\theta$  has a simple zero at  $z_0$ , the first logarithmic derivative has a simple pole at  $z_0$ , and then the rest follows by taking higher derivatives.  $\square$

We also need some more complex analysis

**Lemma 2.8.** Any  $\Lambda_\tau$ -periodic meromorphic function that, up to  $\Lambda_\tau$ -translations, has at most a simple pole is constant.

*Proof.* Suppose that  $f$  is a meromorphic function such that  $f(z + m + \tau n) = f(z)$  for all  $m, n \in \mathbb{Z}$ . Assume that, up to translation,  $f$  has at most one simple pole  $x$ . We can put  $x$  into a fundamental parallelogram  $P$  and then the Residue Theorem tells us that

$$\text{Res}(f, x) = \frac{1}{2\pi i} \int_{\partial P} f(z) dz$$

Since  $f$  is  $\Lambda_\tau$ -periodic, the integral on the right is zero, so that the function has no residues at  $f$ . But this means that  $x$  is not a pole of  $f$ . Thus,  $f$  must actually be holomorphic. Since  $f: \mathbb{C} \rightarrow \mathbb{C}$  is  $\Lambda_\tau$ -invariant, it factors through a holomorphic function  $\bar{f}: \mathbb{C}/\Lambda_\tau \rightarrow \mathbb{C}$  but since the space  $\mathbb{C}/\Lambda_\tau$  is a complex torus, it is compact, and its image is bounded. Then  $f$  is a holomorphic bounded function, hence it is constant.  $\square$

With this we can give the result that we were looking for:

**Theorem 2.9.** *There exist  $a, b, c \in \mathbb{C}$  such that*

$$\left(\frac{\partial^3 \log \theta}{\partial z^3}\right)^2 = -4 \left(\frac{\partial^2 \log \theta}{\partial z^2}\right)^3 + 4a \cdot \left(\frac{\partial^2 \log \theta}{\partial z^2}\right)^2 + b \cdot \left(\frac{\partial^2 \log \theta}{\partial z^2}\right) + c$$

*Proof.* A computation, for example on a computer, shows that for an arbitrary function  $f(z)$  it holds that

$$\begin{aligned} \left(\frac{\partial^3 \log f}{\partial z^3}\right)^2 + 4 \left(\frac{\partial^2 \log f}{\partial z^2}\right)^3 &= \frac{4f'''(z)f'(z)^3}{f(z)^4} - \frac{3f'(z)^2f''(z)^2}{f(z)^4} \\ &+ \frac{4f''(z)^3}{f(z)^3} - \frac{6f'(z)f''(z)f'''(z)}{f(z)^3} + \frac{f'''(z)^2}{f(z)^2} \end{aligned}$$

The precise form is not important, but what is important is the highest power of  $f(z)$  appearing in a denominator is 4. Hence, if we apply this to the theta function, we see that the expression

$$\left(\frac{\partial^3 \log \theta}{\partial z^3}\right)^2 + 4 \left(\frac{\partial^2 \log \theta}{\partial z^2}\right)^3$$

can be written as the quotient of an holomorphic function by  $\theta(z, \tau)^4$ . Hence, the full expression has a pole of order at most 4 at  $z_0 = \frac{1}{2} + \frac{1}{2}\tau$ . We can subtract an appropriate multiple of  $\left(\frac{\partial^2 \log \theta}{\log z^2}\right)^2$  and obtain a function with a pole at  $z_0$  of order at most 3, we can then subtract an appropriate multiple of  $\frac{\partial^3 \log \theta}{\log z^3}$  to obtain a function with a pole of order at most 2 and then we can subtract an appropriate multiple of  $\frac{\partial^2 \log \theta}{\log z^2}$  to obtain a function with a pole of order at most 1. To, summarize, there are  $a, b, d \in \mathbb{C}$  such that the  $\Lambda_\tau$ -periodic function

$$\left(\frac{\partial^3 \log \theta}{\partial z^3}\right)^2 + 4 \left(\frac{\partial^2 \log \theta}{\partial z^2}\right)^3 - 4a \cdot \left(\frac{\partial^2 \log \theta}{\partial z^2}\right)^2 - d \cdot \left(\frac{\partial^3 \log \theta}{\partial z^3}\right) - b \cdot \left(\frac{\partial^2 \log \theta}{\partial z^2}\right)$$

has a pole of order at most 1 at  $z_0$  and, up to  $\Lambda_\tau$ -translation, nowhere else. Then Lemma 2.8 shows that the function must be constant. Hence there is  $c \in \mathbb{C}$  such that

$$\left(\frac{\partial^3 \log \theta}{\partial z^3}\right)^2 = -4 \left(\frac{\partial^2 \log \theta}{\partial z^2}\right)^3 + 4a \cdot \left(\frac{\partial^2 \log \theta}{\partial z^2}\right)^2 - d \cdot \left(\frac{\partial^3 \log \theta}{\partial z^3}\right) + b \cdot \left(\frac{\partial^2 \log \theta}{\partial z^2}\right) + c$$

To conclude, we just need to show that  $d = 0$ . But this is true because the function  $\frac{\partial^3 \log \theta}{\partial z^3}$  is odd, while all other functions appearing are even.  $\square$

This way, we get a solution to the KdV equation

**Corollary 2.10.** *With the same notation as in Theorem 2.9, the function*

$$u(x, t) = 2 \cdot \frac{\partial^2 \log \theta}{\partial z^2}(x + a \cdot t)$$

*is a solution to the KdV equation*

These solutions are usually called *quasiperiodic*.

### 3 Degenerations and soliton solutions

We can get other solutions to the KdV by degenerating the theta function:

**Lemma 3.1.** *Let  $t \in \mathbb{R}$  be a positive real number. Then*

$$\lim_{t \rightarrow +\infty} \theta \left( z - \frac{1}{2}it, it \right) = 1 + \exp(2\pi iz)$$

*Proof.* We just write everything explicitly:

$$\begin{aligned} \theta \left( z - \frac{1}{2}it, it \right) &= \sum_{n \in \mathbb{Z}} \exp \left( \pi in^2 \cdot it + 2\pi in \left( z - \frac{1}{2}it \right) \right) \\ &= \sum_{n \in \mathbb{Z}} \exp \left( -\pi(n^2 - n)t \right) \cdot \exp(2\pi inz). \end{aligned}$$

We see that  $n^2 - n \geq 0$  for all  $n \in \mathbb{Z}$  and when  $n^2 - n > 0$  the term  $\exp(-\pi(n^2 - n)t) \cdot \exp(2\pi inz)$  goes to zero as  $t \rightarrow +\infty$ . The only terms surviving are those for which  $n^2 - n = 0$ : this means precisely  $n = 0, 1$  so that the limit is exactly what we want.  $\square$

This way we have obtained a sort of degenerate theta function

$$\hat{\theta}(z) = 1 + \exp(2\pi iz)$$

Do we still get a solution of the KdV equation out of this? One can compute that

$$\left( \frac{\partial^3 \log \hat{\theta}}{\partial z^3} \right)^2 + 4 \left( \frac{\partial^2 \log \hat{\theta}}{\partial z^2} \right)^3 = (2\pi i)^2 \cdot \left( \frac{\partial^2 \log \hat{\theta}}{\partial z^2} \right)^2$$

so that the function

$$u(x, t) = 2 \frac{\partial^2 \log \hat{\theta}}{\partial z^2} \hat{\theta}(x + \pi^2 t)$$

is a solution to the KdV equation thanks to Lemma 1.1. This is what one could call a soliton solution to the KdV equation. Geometrically, what is happening is that the curve  $y^2 = -4x^3 - 4\pi^2 \cdot x^2$  is now a singular cubic curve, with a nodal singularity at  $(x, y) = (0, 0)$ .