

Born's rule A self-adj. op. in \mathcal{H} .

$$\rho_A(\alpha) = \sum_{\lambda} |\langle \phi_{\alpha, \lambda} | \psi \rangle|^2 = \sum_{\lambda} |(\mathcal{U}\psi)(\alpha, \lambda)|^2$$

Conseq $E(\alpha) = \int d\alpha \alpha \rho_A(\alpha)$ $\mathcal{U}: \mathcal{H} \rightarrow L^2(\Omega)$

$$= \int_{\Omega} d(\alpha, \lambda) \alpha \cancel{\sum_{\lambda}} |(\mathcal{U}\psi)(\alpha, \lambda)|^2$$

$$\left[\begin{array}{l} M = \text{mult. by} \\ \alpha: \Omega \rightarrow \mathbb{R} \\ M: L^2(\Omega) \rightarrow L^2(\Omega) \\ \text{or dense subspace therein} \end{array} \right] = \langle (\mathcal{U}\psi | (\mathcal{U}\mathcal{U}^{-1}M \mathcal{U}\psi) \rangle_{\Omega}$$

$$= \langle \psi | \underbrace{\mathcal{U}^{-1}M\mathcal{U}}_A \psi \rangle_{\mathcal{H}} = \langle \psi | A \psi \rangle$$

Apply spectral theorem to H

$$H \phi_{E,\lambda} = E \phi_{E,\lambda},$$

$\{\phi_{E,\lambda}\}$ GONB
energy eigenbasis

$$\psi = \sum_{E,\lambda} c_{E,\lambda} \phi_{E,\lambda}$$

$$\psi_t = \underbrace{e^{-iHt/\hbar}} \psi = \sum_{E,\lambda} \underbrace{e^{-iEt/\hbar}} c_{E,\lambda} \phi_{E,\lambda}$$

i.e., $c_{E,\lambda}(t) = e^{-iEt/\hbar} c_{E,\lambda}(0)$

Conservation laws in QM

$E(t), \underline{p}(t), \underline{L}(t)$ t -indep. in Newtonian mech.
no analog in QM.

$|c_{E,\lambda}(t)|^2$ are t -indep. : conservation law,
no analog in Newtonian mech.

H, P, L op.s are conserved:

ONB $\{\phi_i\}$, $\{\phi_i(t)\}$ ONB, $S_{ij}^{(t)} = \langle \phi_i^{(t)} | S | \phi_j^{(t)} \rangle$ t -indep.

if $\underline{e^{iHt} S e^{-iHt} = S} \iff \underline{SH = HS}$

"S is conserved"

Conseq for actual values

$$PH_0 = H_0 P$$

$$P H \neq H P$$

1) quantum measurement of S
no conservation of S .

2) Ex tunneling, $H = H_0 + V$

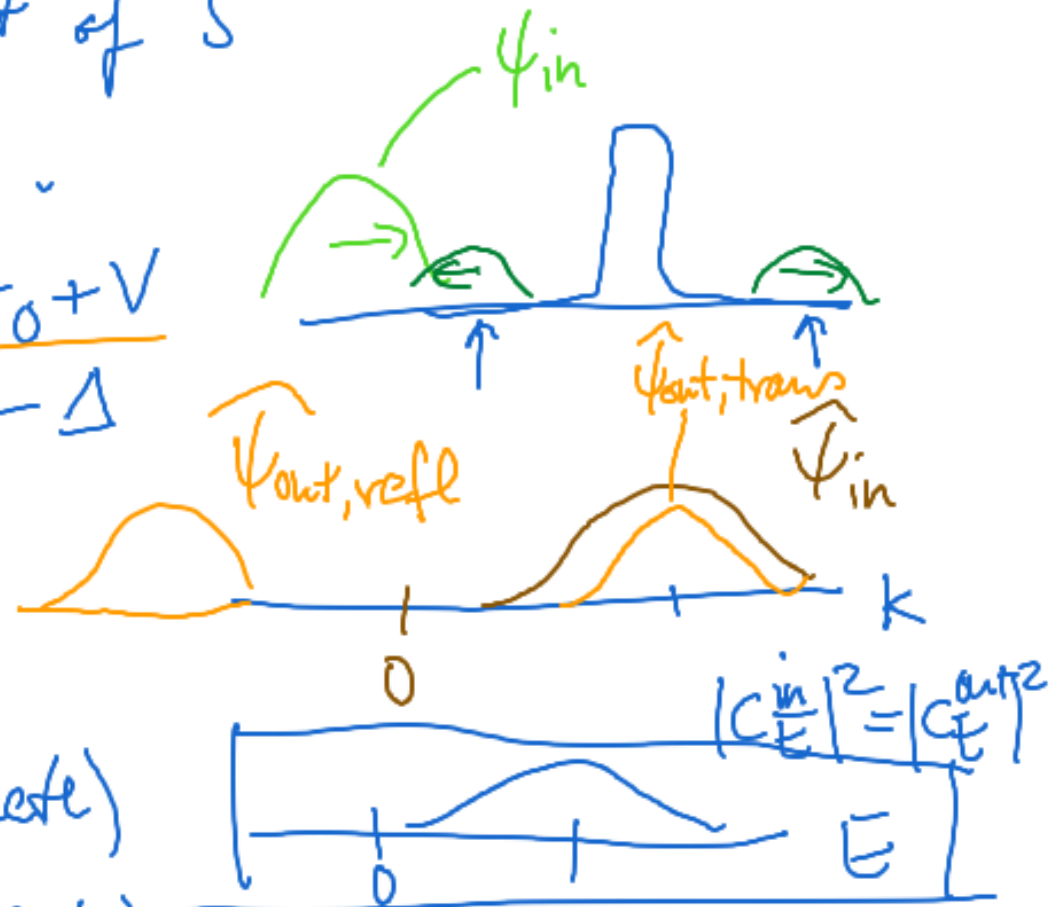
early: $\rho_{PH}(E) \approx \rho_{H_0}(E)$ $H_0 = -\Delta$

late: $\rho_H(E) \approx \rho_{H_0}(E)$

$\rho_H(E, t)$ is t -indep.

$$\rho_H(E, \text{early}) = \rho_H(E, \text{late})$$

$$\rho_{H_0}(E, \text{early}) \approx \rho_{H_0}(E, \text{late})$$



Spin

$$\psi: \mathbb{R}^3 \rightarrow \mathbb{C} \quad \text{before}$$

$$\psi: \underline{\mathbb{R}^3} \rightarrow \underline{\mathbb{C}^2} \quad \text{now}$$

$$\mathbb{C}^2 = \text{spin space} = S$$

$$\cancel{S} \ni \phi \quad \text{spin vector} \\ \text{spinor}$$

Spinors

S is a Hilbert space, $\dim S = 2$.

$S \ni \phi$ is associated with a vector in physical space \mathbb{R}^3 .

$$\underline{\omega}(\phi) = \left(\sum_{r,s=1}^2 \phi_r^* (\sigma_1)_{rs} \phi_s, \dots, \sigma_2 \dots, \dots, \sigma_3 \dots \right)$$

$\omega_a(\phi) = \langle \phi | \sigma_a \phi \rangle_S$

$$\langle \phi | \chi \rangle = \sum_{s=1}^2 \phi_s^* \chi_s = \phi^* \chi = \phi^\dagger \chi$$

Pauli matrices: $\underline{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ +i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

self-adj, complex 2×2 , have e.v.s ± 1 ,

$$\underline{\text{Ex}} \quad \phi = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \underline{\omega}(\phi) = \left(\phi^* \sigma_1 \phi, \phi^* \sigma_2 \phi, \phi^* \sigma_3 \phi \right)$$

$$\phi^* M \phi = M_{11} = (0, 0, 1)$$

"up spinor"

$$\phi = \begin{pmatrix} 0 \\ 1 \end{pmatrix}, \quad \underline{\omega}(\phi) = (0, 0, -1) \quad \text{"down spinor"}$$

$$\phi^* M \phi = M_{22}$$

Properties • $\underline{\omega}(z\phi) = |z|^2 \underline{\omega}(\phi) \quad z \in \mathbb{C}$

• $\|\underline{\omega}(\phi)\|_{\mathbb{R}^3} = \|\phi\|_S^2 = \phi^* \phi \quad (\text{exercise})$

• Def angle in \mathcal{H} : $\cos \theta = \frac{|\langle \phi | \chi \rangle|}{\|\phi\| \|\chi\|}$

Thm Suppose $\phi, \chi \in S$, $\|\phi\|_S = 1 = \|\chi\|_S$ with angle θ . Then $\text{angle}_{\mathbb{R}^3}(\underline{\omega}(\phi), \underline{\omega}(\chi)) = 2\theta$.

Spin- $\frac{1}{2}$. Ex $\begin{pmatrix} 1 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ have $\theta = 90^\circ$
 $\underline{\omega}\left(\begin{pmatrix} 1 \\ 0 \end{pmatrix}\right) = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}, \underline{\omega}\left(\begin{pmatrix} 0 \\ 1 \end{pmatrix}\right) = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$ have angle 180° .

Other spins

spinors of $\text{spin}-\frac{1}{2}$, $\text{spin}-1$, $\text{spin}-\frac{3}{2}$, $\text{spin} 2$, $\text{spin}\frac{5}{2}$, ...

The space of $\text{spin}-s$ spinors has $\dim 2s+1$.

$\text{spin}-0$: $\dim=1$, $\psi: \mathbb{R}^3 \rightarrow \mathbb{C}$

$\text{spin}-s$:

$$(\sigma_1)_{rr'} = \frac{1}{2} (\delta_{r,r'+1} + \delta_{r+1,r'}) \sqrt{(2s+1)(r+r'-1) - rr'}, \quad r, r' \in \{1 \dots 2s+1\}$$

$$(\sigma_2)_{rr'} = \frac{-i}{2} (\delta_{r,r'+1} - \delta_{r+1,r'}) \sqrt{(2s+1)(r+r'-1) - rr'}$$

$$(\sigma_3)_{rr'} = \delta_{rr'} (2s+1 - r)$$



in nature, electrons have spin $\frac{1}{2}$
quarks have spin $\frac{1}{2}$

all matter particles have spin $\frac{1}{2}$.

force particles have integer spins
either 1 or 2

photons 1

gluons 1

gravitons 2

except Higgs boson 0

(confirmed in 2012)