

Spin: Actual spin values vs. Actual Positions

$$L^2(\mathbb{R}^3, \mathbb{C}^2) \cong L^2(\underline{\mathbb{R}^3 \times \{1,2\}}, \mathbb{C})$$

$$\psi(x) \in \mathbb{C}^2 \quad \psi(x, s)$$

may suggest  $(Q, \bar{s})$  Bohm-style theory

$$\text{Prob}(Q \in d^3q, \bar{s} = s) = |\psi(q, s)|^2 d^3q$$

Problem: not rotationally invariant.

Upshot: resist temptation.

## Representations of $SO(3)$

$\psi: \mathbb{R}^3 \rightarrow \mathbb{C}^3$ , change of  $\psi$   $R \in SO(3)$

$$\tilde{\psi}(x) = R \psi(R'x)$$

tensor-valued:  $\psi: \mathbb{R}^3 \rightarrow \mathbb{C}^3 \otimes \mathbb{C}^3$

$$\psi_{ab}, \quad \tilde{\psi}_{ab}(x) = \sum_{c,d=1}^3 R_{ac} R_{bd} \psi_{cd}(R'x)$$

pattern:  $\psi: \mathbb{R}^3 \rightarrow \mathbb{C}^d$ ,  $\tilde{\psi}_r(x) = \sum_{s=1}^d M_{rs}(R) \psi_s(R'x)$

composition law:  $M(R_1) M(R_2) = M(R_1 R_2)$ , transformation law:  $M(I) = I$

gr. rep : gr. homomorphism  $SO(3) \rightarrow \underline{GL}(\mathbb{C}^d)$

$S'$  for spin- $\frac{1}{2}$ , one more complication:

$\forall R \in SO(3)$ . not  $S' \rightarrow S$  but  $\underline{P}(S') \rightarrow P(S)$

projective space  $P(V) = \{1\text{d subspaces of } V\}$ .

$e^{i\theta} \psi$  equal to  $\psi$  in Born rule  $\forall \theta$   
in BM

"ray" = 1d subsp.

projective mapping  $F: P(\mathcal{H}) \rightarrow P(\mathcal{H})$

compatible with a linear mapping  $M: \mathcal{H} \rightarrow \mathcal{H}$   
 $F(C\psi) = CM\psi$  ( $\ker M = \{0\}$ )

projective repr.  $F(R_1) F(R_2) = F(R_1 R_2)$ ,  $F(I) = \bar{I}$

linear

$$\underline{M(R_1)} M(R_2) = \pm \underline{M(R_1 R_2)}$$

gr. homom.  $SO(3) \rightarrow$  projective mappings( $S$ )

$M \in GL(\mathbb{C}^2)$ ,  $M \in SU(2)$ ,

$\varphi: SU(2) \rightarrow SO(3)$  gr. homom.

$\varphi$  smooth, 2-to-1, locally diffeomorphism

Ex  $x: \mathbb{R} \rightarrow U(1)$ ,  $\theta \mapsto e^{i\theta}$  is gr. homom., smooth, many-to-one, locally diffeom.

unfold  $U(1)$ , get  $\mathbb{R}$

unfold  $SO(3)$ , get  $Sl(2)$

"unfold" = universal covering space

Curve in  $SO(3)$  starting from  $\overline{I}$

$\exists$  curve in  $SU(2)$  starting from  $\overline{I}$   
called the lift



$SO(3)$



$SU(2)$

2 ways:  $\xrightarrow[\text{gr. repr.}]{\text{proj.}} SO(3) \rightarrow P(\mathbb{C}^2)$

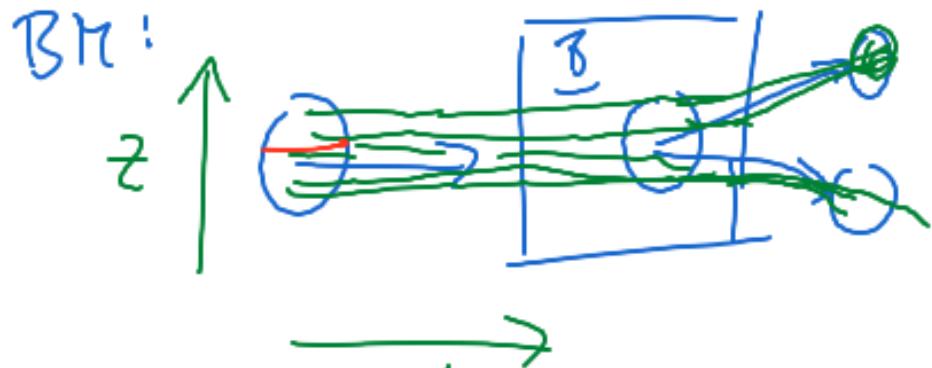
or  $\xrightarrow[\text{gr. repr.}]{\text{proj.}} SU(2) \rightarrow GL(\mathbb{C}^d)$

$S$  is equipped with a particular  
repr. of  $\widehat{SO(3)} \cong SU(2)$  / proj. repr. of  $SO(3)$ .

Def  $S$  is the repr. space of  $\widehat{SO(3)}$   
of the "spin- $\frac{1}{2}$ " repr.

## Inverted SG magnet and contextuality

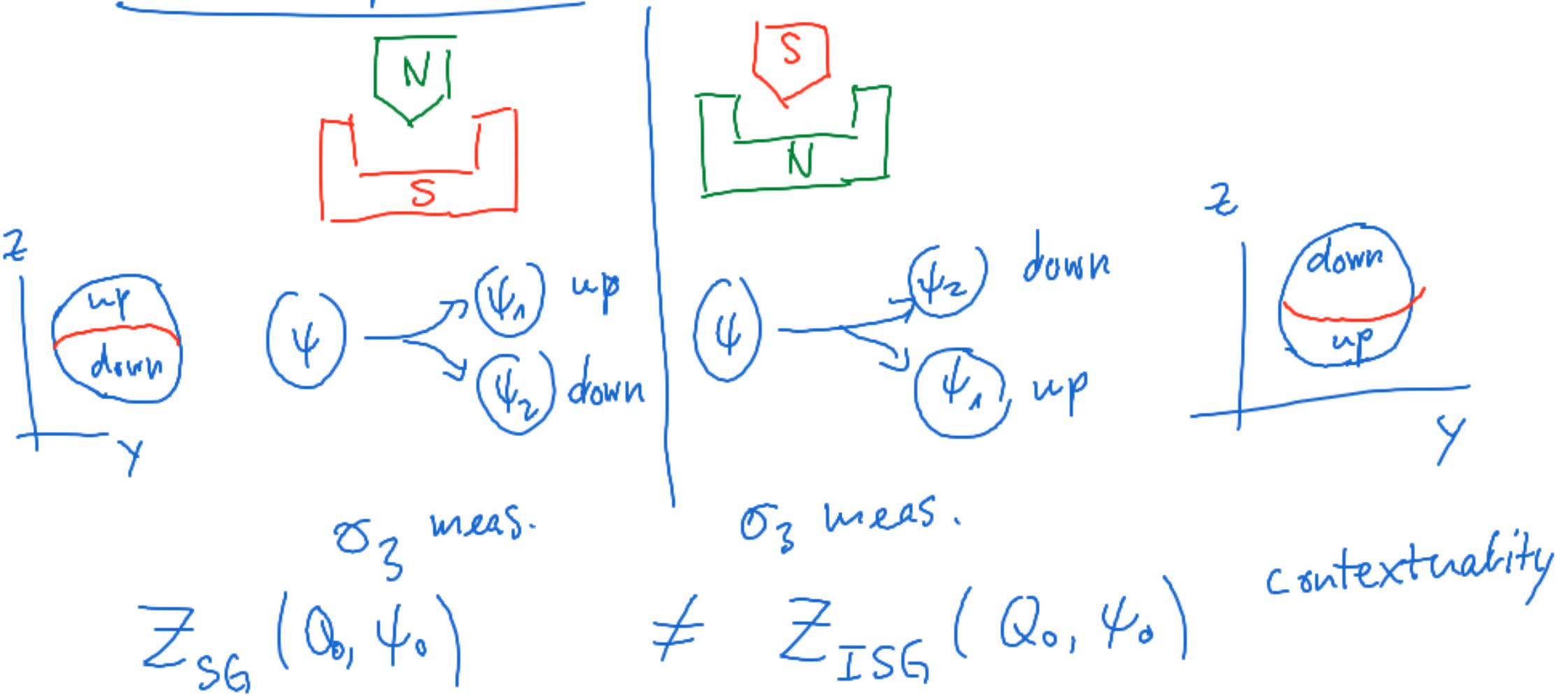
SG in  direction,  $\psi_s(x) = \underline{\phi_s} \chi(x)$



The result of the experiment is determined by  $\underline{Q}(0)$  and  $\psi(0)$

In general, the result is not determined by  $\underline{Q}(0)$ ,  $\psi(0)$ , and  $\underline{\phi}_3$ .

## Another experiment



# Collapse and Measurement

## The Projection Postulate

Notation: Dirac  $|\psi\rangle = \psi \in \mathcal{H}$

$$|\phi_n\rangle = |n\rangle \text{ ket}$$

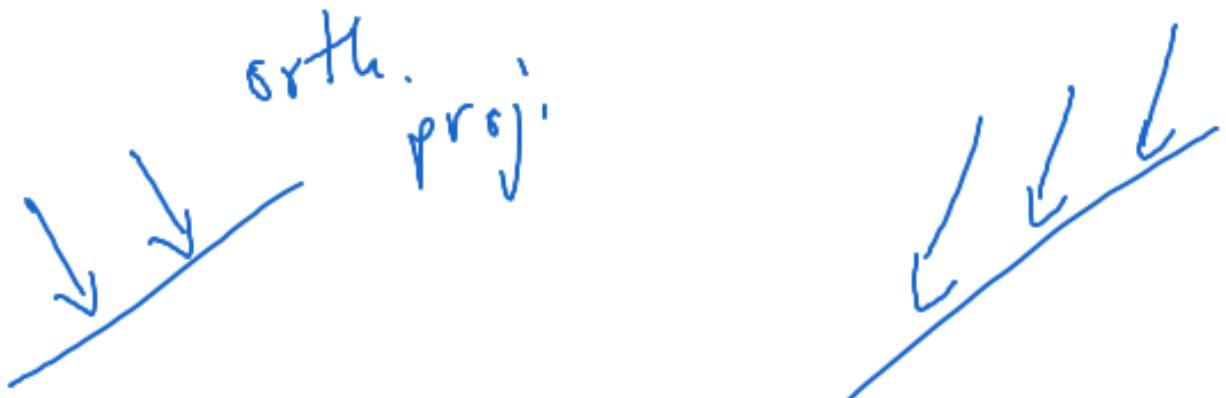
$$\langle\phi| : \mathcal{H} \rightarrow \mathbb{C} : \psi \mapsto \langle\phi|\psi\rangle$$

bra  $\langle\phi| |\psi\rangle = \langle\phi|\psi\rangle$

$$\langle\phi| A\psi\rangle = \underbrace{\langle\phi| A|\psi\rangle}_{\text{particularly for self-adj. } A.}$$

$$|\phi\rangle\langle\phi| : \mathcal{H} \rightarrow \mathcal{H}$$

$|\phi\rangle\langle\phi|$ , if  $\|\phi\|=1$ , then this is the proj. to  $\mathbb{C}\phi$ .



⊗ notation:  $\underline{\Psi(x,y)} = \psi(x) \phi(y)$   
 $\Leftrightarrow \Psi = \psi \otimes \phi$

$$L^2(x,y) = L^2(x) \otimes L^2(y) \quad \text{tensor product}$$

$$\mathbb{C}^2 \otimes L^2(\mathbb{R}^3, \mathbb{C}) = L^2(\mathbb{R}^3, \mathbb{C}^2)$$

Notation:  $f(t-) = \lim_{s \rightarrow t^-} f(s)$ ,  $f(t+) = \lim_{s \searrow t^+} f(s)$