

$$\underline{\text{Thm}} \quad A = A^\dagger, \quad B = B^\dagger, \quad \| \psi \| = 1, \quad \sigma_A \sigma_B \geq \frac{1}{2} | \langle \psi | [A, B] | \psi \rangle |$$

? f Born distr _{ψ, A} , $\langle A \rangle = \langle \psi | A \psi \rangle$ expectation

$$\sigma_A^2 = \langle \psi | (A - \langle A \rangle)^2 | \psi \rangle = \| \phi_A \|^2, \quad \phi_A := (A - \langle A \rangle) \downarrow \psi$$

Cauchy-Schwarz inequality

$$| \langle \phi_A | \phi_B \rangle |^2 \leq \| \phi_A \|^2 \| \phi_B \|^2 = \sigma_A^2 \sigma_B^2$$

$$\begin{aligned} \text{Now } \langle \phi_A | \phi_B \rangle &= \langle \psi | (A - \langle A \rangle) (B - \langle B \rangle) | \psi \rangle \\ &= \langle \psi | AB - \underbrace{\langle A \rangle B - A \langle B \rangle + \langle A \rangle \langle B \rangle}_{\text{independent}} | \psi \rangle \\ &= \langle \psi | AB | \psi \rangle - \langle A \rangle \langle B \rangle - \cancel{\langle A \rangle \langle B \rangle} + \cancel{\langle A \rangle \langle B \rangle} \\ &=: \underline{\langle AB \rangle - \langle A \rangle \langle B \rangle} \end{aligned}$$

$$\underbrace{|\langle \phi_A | \phi_B \rangle|^2}_{\geq} \geq (\Im m \langle \phi_A | \phi_B \rangle)^2$$

$$= \left| \frac{\langle \phi_A | \phi_B \rangle - \langle \phi_B | \phi_A \rangle}{2i\lambda} \right|^2$$

$$= \frac{1}{4} \left| \langle AB \rangle - \cancel{\langle A \rangle \langle B \rangle} - \cancel{\langle BA \rangle} + \cancel{\langle B \rangle \langle A \rangle} \right|^2$$

$$= \frac{1}{4} \left| \langle \Psi | [A, B] | \Psi \rangle \right|^2$$

□

The Measurement Problem

Consider quantum meas. of $A = A^\dagger$, spectrum (A) is finite
object + apparatus "ready state"

$$\Psi(x,y) = \psi(x) \quad \overset{\phi(y)}{\overbrace{\psi(y)}}, \quad \Psi(t) = \psi \otimes \phi \text{ at } t_1 \text{ (begin)}$$
$$\Psi(t_2) \quad \quad \quad t_2 \text{ (ends)}$$

$$A = \sum_{\alpha} \alpha P_{\alpha}, \quad \psi = \sum_{\alpha} c_{\alpha} \psi_{\alpha}, \quad \|\psi_{\alpha}\| = 1, \quad \sum_{\alpha} |c_{\alpha}|^2 = 1$$

First, suppose $\psi = \psi_{\alpha}$. Born $\Rightarrow P(Z = \alpha') = \delta_{\alpha \alpha'}$.

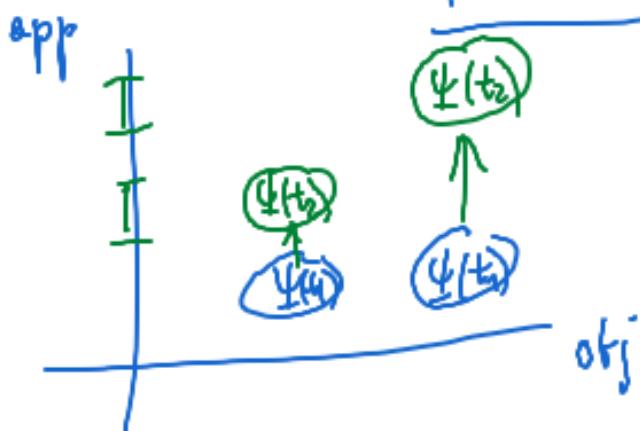
Let $\Psi_{\alpha}(t_1) = \psi_{\alpha} \otimes \phi$. Evolves to $\Psi_{\alpha}(t_2)$ (displays α)

Second, general $\psi = \sum c_\alpha \psi_\alpha$

$$\Psi(t_1) = \psi \otimes \phi = \sum c_\alpha \Psi_\alpha(t_1)$$

Schr. \downarrow

$$\boxed{\Psi(t_2) = \sum c_\alpha \Psi_\alpha(t_2)}$$

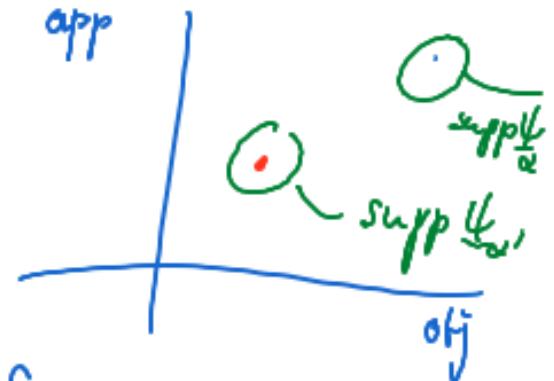


- 1) In each run, there is a unique outcome.
- 2) The wavefct is the complete description of a system's physical state.
- 3) Time evol. of isolated system is linear.

- 1) is dropped in the many-worlds theory
- 2) is dropped in B.H.
- 3) is dropped in GRW.

How BM solves the meas. problem

$$\Psi(t_2) = \sum_{\alpha} c_{\alpha} \Psi_{\alpha}$$



$$P(Q \in \text{support}(\Psi_{\alpha})) = \int_{\text{support}(\Psi_{\alpha})} dq |\Psi_{\alpha}(q)|^2 = \int_{\mathbb{R}^{3N}} dq |c_{\alpha} \Psi_{\alpha}(q)|^2$$

$$= |c_{\alpha}|^2 \|\Psi_{\alpha}\|^2 = |c_{\alpha}|^2 \|\Psi_{\alpha}(t_1)\|^2 = |c_{\alpha}|^2 \|\Psi_{\alpha} \otimes \phi\|^2$$

$= |c_{\alpha}|^2$ in agreement with Born's rule for A.

Conseq BM agrees with quantum formalism
No experiment can distinguish between BM & QF.

Collapse of wf in BM

① Ψ_α will never overlap again ("decoherence")
not for $10^{10^{10}}$ yrs.

b/c they have macroscopically disjoint supports.



$$\Psi = \sum c_\alpha \Psi_\alpha \rightsquigarrow \Psi_{\alpha_0}$$

② wf of object

If $\Psi(x,y) = \psi(x)\phi(y)$ then Ψ .

general Def conditional wf

$$\Psi_{\text{cond}}(x) = \frac{1}{N} \int dY |\Psi(x,Y)|^2$$

doesn't evolve according to Schr. eq.,

except when $\Psi(x,y) = \psi(x)\phi(y)$, and no interaction

$$\Psi(t_2) = \sum c_j \Psi_j, \text{ suppose } \Psi_j = \psi_j \otimes \phi_j$$