

Probability in Many-Worlds Theories

1000 SG, $|c_{up}|^2 = \frac{1}{3}$, $|c_{down}|^2 = \frac{2}{3}$.

↑ ↓ ↓ ↑ ↓ ↑ ↓ ↓ ...

Rule for counting worlds. (for SG)



The "fraction of worlds" $f(P)$ with property P in the splitting

given by $\Psi = \sum c_\alpha \Psi_\alpha$ and $m(x) = \sum_\alpha |c_\alpha|^2 m_\alpha(x)$ is

$$f(P) = \sum_{\alpha \in M} |c_\alpha|^2$$

where M is the set of worlds α with property P .

Ex $P =$ freq. of \uparrow lies in $[30\%, 36\%]$

Comparison to BM: $\Psi_0, Q_0 \sim |\Psi_0|^2$

$\Rightarrow Q_t \sim |\Psi_t|^2$, P true for some Q_t
false for some Q_t .

typicality reasoning.

(Everett 1955)

Robustness against re-drawing boundaries between branches



$\|P_\alpha \Psi\|^2$



Problem with the rule for counting worlds

Is there room for such a rule among the laws of physics?

I tend to conclude: in the end, S_{MW} fails.

Comparison: Bell's MW #1: OK.

Bell's MW #2. $Q_t \sim |\Psi_t|^2$, indep. of past.
OK as well.

Other reasonings for justifying probabilities

D. Deutsch (1999): rational behavior

L. Vaidman (1998): I can be ignorant of which world I am in.

Nonlocality

The Einstein-Podolsky-Rosen Argument

EPR 1935 argued for

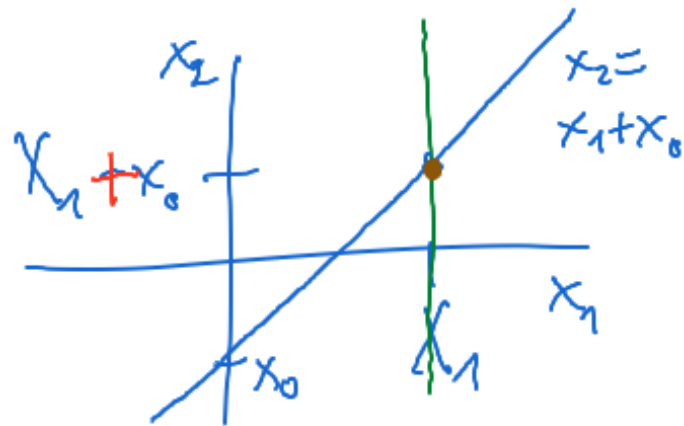
claim: There are additional variables beyond the wave fct.

"PF" 2 particles in 1d with wf
$$\Psi(x_1, x_2) = \delta(x_1 - x_2 + x_0)$$

$x_0 \in \mathbb{R}$ const.

Alice measures of x_1 , outcome $X_1 \sim$ uniform.

$$\Psi'(x_1, x_2) = \delta(x_1 - X_1) \Psi(x_1, x_2) = \delta(x_1 - X_1) \delta(x_2 - X_1 - x_0)$$



Now Bob measures x_2 , obtains $X_2 = X_1 + x_0$
with certainty

\Rightarrow before Bob's measurement, particle 2 had a definite position.

EPR assumption: "no real change can take place in the 2nd system in conseq. of [a measurement on] the first system."

(a special case of locality)

\Rightarrow before Alice's measurement, particle 2 had a definite position. \square



Variant: squ.-integrale



$$\underbrace{\begin{pmatrix} P & \otimes I \\ B \end{pmatrix}}_{P_{B \times R}}$$