

Bogoliubov maps / Bogoliubov transformations

We have learned that for any creation op $a^*(t)$ it holds:

$$I: U_{t,s}^{Bog} a^*(t) U_{s,t}^{Bog} = a(\eta) + a^*(\rho)$$

resp. $II: U_{t,s}^{Bog} a(t) U_{s,t}^{Bog} = a(\tilde{\eta}) + a^*(\tilde{\rho})$

Properties: a) The equations must still hold taking the adjoint:

$$\Rightarrow I^{adj}: U_{t,s}^{Bog} a(t) U_{s,t}^{Bog} = a^*(\eta) + a(\rho)$$

If we summarize I and II in matrix form:

$$\begin{pmatrix} f \\ g \end{pmatrix} \rightarrow \mathbb{T} \begin{pmatrix} f \\ g \end{pmatrix} = \begin{pmatrix} T_{11} & T_{12} \\ T_{21} & T_{22} \end{pmatrix} \begin{pmatrix} f \\ g \end{pmatrix} = \begin{pmatrix} T_{11}f + T_{12}g \\ T_{21}f + T_{22}g \end{pmatrix}$$

$$U_{t,s}^{Bog} a(t) U_{t,s}^{Bog} = a(T_{11}(t)) + a^*(T_{21}(t))$$

$$U_{t,s}^{Bog} a^*(g) U_{t,s}^{Bog} = a(T_{12}(g)) + a^*(T_{22}(g))$$

$$\Rightarrow \begin{cases} T_{12} = \overline{T_{21}} \\ T_{11} = \overline{T_{22}} \end{cases}$$

$$\mathbb{T} = \begin{pmatrix} U & W \\ \bar{W} & \bar{U} \end{pmatrix}$$

b) Further $[a(t), a^*(t)] = 1$ for t normalized.

$$U_{t,s}^{Bog} [a(t), a^*(t)] U_{s,t}^{Bog} = 1$$

$$[a(U(t)) + a^*(\bar{W}(t)), a(W(t)) + a^*(\bar{U}(t))] = 1$$

$$UU^* - W^*W = 1$$

$$\langle U(t), \bar{U}(t) \rangle - \langle W(t), \bar{W}(t) \rangle = 1$$

Also $[a(t), a(t)] = 0$

$$\Rightarrow [a(U(t)) + a^*(\bar{W}(t)), a(U(t)) + a^*(\bar{W}(t))] = 0$$

$$\langle U(t), \bar{W}(t) \rangle - \langle U(t), \bar{W}(t) \rangle = 0 \quad \text{etc.}$$

Definition: A linear map of the form $\begin{pmatrix} U & W \\ \bar{W} & \bar{U} \end{pmatrix}$ that satisfies the properties induced by the commutation relations is called a Bogoliubov map.

In case there is a unitary transformation S such that

$$S a S^{-1} = a(U(t)) + a^*(\bar{W}(t)) \quad \text{etc.}$$

is called Bogoliubov transformation.

Theorem: If $\text{Tr}(W^*W)$ is finite, then there is a Bog. trafo for

$$\begin{pmatrix} U & W \\ \tilde{W} & \tilde{U} \end{pmatrix}.$$

Beyond Bog. dynamic

Using some perturbation argument one can describe the free time evolution with many details. We do so using Duhamel expansion:

$$\| U_{t,s} \psi - \tilde{U}_{t,s} \psi \| = \int_s^t d\tau \tilde{U}_{t,\tau} U_{\tau,s} \psi \stackrel{\leq}{=} \int_s^t \tilde{U}_{t,\tau} \| (\tilde{H} - H) U_{\tau,s} \psi \| d\tau$$

Duhamel expansion:

$$\rightarrow \tilde{\psi}_t = \tilde{U}_{t,0} \psi = U_{t,0} \psi - i \int_0^t \tilde{U}_{t,s} (\tilde{H} - H) U_{s,0} \psi ds - \dots$$

$[\tilde{U}_{t,\tau} U_{\tau,s} \psi]_s = \tilde{U}_{t,t} U_{t,s} \psi - \tilde{U}_{t,s} U_{s,s} \psi$

$$\langle \tilde{\psi}_t, M \tilde{\psi}_t \rangle = \langle \psi_t, M \psi_t \rangle + \langle \psi_t, M (-i \int_0^t U_{t,s} (\tilde{H} - H) U_{s,0} \psi_0 ds) \rangle_{0,0}$$

$$\langle \underbrace{R_1 U_{t,0}^{\text{Bog}} M U_{0,t} R_2}_{T(M)} \rangle + \langle \underbrace{\psi_t}_{U_{t,0}} \uparrow, M (-i \int_0^t \underbrace{U_{t,s}^{\text{Bog}}}_{U_{s,t}} (\tilde{H} - H) U_{s,0}^{\text{Bog}} \psi_0 ds) \rangle_{0,0}$$