# **SCARS ON GRAPHS**

#### Holger Schanz and Tsampikos Kottos





Nichtlineare Dynamik Universität Göttingen

Max-Planck-Institut

für Strömungsforschung Göttingen

### Deterministic Chaos vs Random Walk on a Graph



1. Line

$$-\frac{\hbar^2}{2m}\frac{\mathrm{d}^2}{\mathrm{d}x^2}\Psi(x)=\mathsf{E}\,\Psi(x)$$

$$k=\sqrt{2mE/\hbar^2}$$

$$a_+ \exp(+ikx)$$

$$- \int a_exp(-ikx)$$

# Quantum graphs:

2. Vertex



Neumann b.c.:

$$\sigma = \begin{pmatrix} \rho & \tau & \tau \\ \tau & \rho & \tau & \cdots \\ \tau & \tau & \rho \\ & \cdots & & \cdots \end{pmatrix}$$

$$egin{array}{rl} au &=& rac{2}{\mathsf{v}} & \ll & 1 & (\mathsf{v} 
ightarrow \infty) \ 
ho &=& rac{2}{\mathsf{v}} - 1 & \sim -1 & (\mathsf{v} 
ightarrow \infty) \end{array}$$

- current conservation  $\sigma\sigma^{\dagger}=\mathbf{I}$
- continuity of wavefunction

# Quantum graphs:

3. Network



interpretation: discrete time evolution  $|\psi(t)\rangle = U^{t}(k)|\psi(0)\rangle$ 

classical analogue: Markov chain  $M_{d'd} = |U_{d'd}|^2$ 

# A complete Graph



# An "ergodic" eigenstate



# A typical eigenstate



### A scar on a graph



#### The inverse participation number



Heller '84: scars  $\Leftrightarrow$  short-time dynamics:

 $\Rightarrow$  average localization of eigenstates

$$\Delta \epsilon \sim \frac{2\pi}{2B} \quad \Rightarrow \quad \langle \cdot \rangle_{t} \sim \frac{1}{2B} \sum_{t=1}^{2B}$$

#### Many ways to return ...



Neumann b.c., large graphs:  $au = 2/v \quad o \quad 0$  $ho = au - 1 \quad o - 1$ 



Kaplan 2001: period-two orbits  $\Rightarrow \langle I \rangle \sim v \times I_{\rm RMT}$ The shortest and most stable orbits cause enhanced localization.

# Take-home message



## Which orbits can scar?



perfect scars:

$$\mathbf{0} = \mathbf{0} + \tau \sum_{\mathsf{d} \in \mathsf{D}_{j,\mathsf{p}}^{(-)}} \mathrm{e}^{\mathrm{i}\mathsf{k}\mathsf{L}_\mathsf{d}} \, \mathsf{a}_\mathsf{d}$$

$$\mathsf{v}_{\mathsf{j},\mathsf{p}} \geq 2 - \delta_{\mathsf{v}_{\mathsf{j}},1} \quad (\forall \mathsf{j} \in \mathsf{p})$$

Stability is irrelevant!



#### Energies of scars?



$$\begin{split} \mathbf{d} &= \sum_{\substack{d' \in \mathsf{D}_{j,p}^{(-)} \\ 0 \\ \end{array}} e^{\mathbf{i}\mathbf{k}\mathsf{L}_{d'}} \mathbf{a}_{d'} \left(\tau + \delta_{d'\hat{d}} \left[\rho - \tau\right]\right) \\ & \underbrace{\mathbf{a}_{d} = -e^{\mathbf{i}\mathbf{k}\mathsf{L}_{d}} \mathbf{a}_{\hat{d}}}_{\mathbf{a}_{d}} = -e^{\mathbf{i}\mathbf{k}\mathsf{L}_{d}} \mathbf{a}_{\hat{d}} \\ & \mathbf{a}_{d} = +e^{2\mathbf{i}\mathbf{k}\mathsf{L}_{d}} \mathbf{a}_{d} \end{split}$$

 $\Rightarrow$  commensurate bond lengths

$$\mathsf{L}_\mathsf{d}/\mathsf{L}_\mathsf{d'} = \mathsf{m}_\mathsf{d}/\mathsf{m}_\mathsf{d'}$$

No perfect scars for generic graphs!

а

#### Perturbation theory for the scar quality



Scar quality:

$$\delta_{\mathsf{p}} = \sum_{\mathsf{d} \notin \mathsf{p}} |\mathsf{a}_\mathsf{d}|^2$$

( $0 \leq \delta \leq 1$ ,  $\delta_p = 0$ : perfect scar)

#### Distribution of scars



$$\mathcal{P}^{(2)}(\delta o 0) \sim \delta^{-1/2} \ \mathcal{P}^{(3)}(\delta o 0) = \mathsf{C} \ \mathcal{P}^{(4)}(\delta o 0) \sim \delta^{+1/2}$$



cf Berkolaiko et al. (2003): No quantum ergodicity for star graphs.

#### Conclusions

- The scar theory of Heller et al. applies to graphs, ...
- ... but it does not describe the scars ...
- ... because strong and weak scarring are unrelated phenomena.
- A detailed understanding of strong scars was achieved, ...
- ... but the method does not (immediately) generalize to other systems.

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